



Chisinau, 2018

Dinamica cuantica a sistemelor opto-mecanice

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• Sumar

- Laboratorul meu
- Punct cuantic in rezonator mechanic
- Punct cuantic in cavitate optica
- Multe puncte cuantice in rezonator mechanic

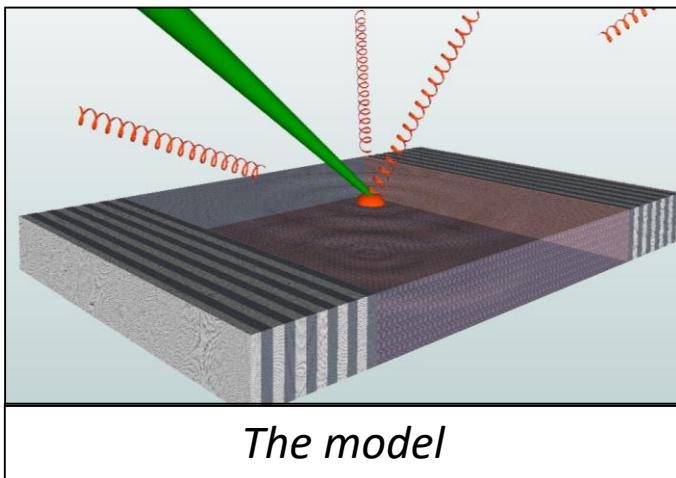
• Laboratorul de Fotonica Cuantica

Studiul teoretic:

- **Opto-mecanica cuantica:** sisteme de nano-rezonatoare mecanice care interactioneaza cu lumina (miscarea materiei la nivel cuantic, determinarea masei moleculelor, racire cuantica)
- **Sisteme de atomi artificiali** care interactioneaza la nivel cuantic cu fotonii (lumina) si fononii (sunet) (interfata cuantica intre electronica si tehnologii cuantice)
- **Disipare cuantica** (pierderea informatiei cuantice, calculatorul cuantic),
corelatiile fotonilor si/sau fononilor in sisteme cuantice
inseparabilitate cuantica (teleportarea semnalului, encriptarea cuantica a informatiei)
- **Cooperare** cu Institutul Max-Planck de fizica nucleara din Heidelberg (MPIK)

Punct cuantic in rezonator mecanic

• Punct cuantic in rezonator mechanic



Cu ce lucram:

- ✓ Un punct cuantic (PC) in rezonator mechanic
- ✓ Laser
- ✓ Temperatura (fononi termali)
- ✓ Intuneric (vacuum electromagnetic)

Ce se intampla

- **Laser albastru defazat => Excitarea PC + crearea fononilor si Dezexcitarea PC + absorbirea fononilor**
- **Laser rosu defazat => Dezexcitarea PC + absorbirea fononilor si Excitarea PC + crearea fononilor**
- **Emisia spontana => doar dezexcitare**

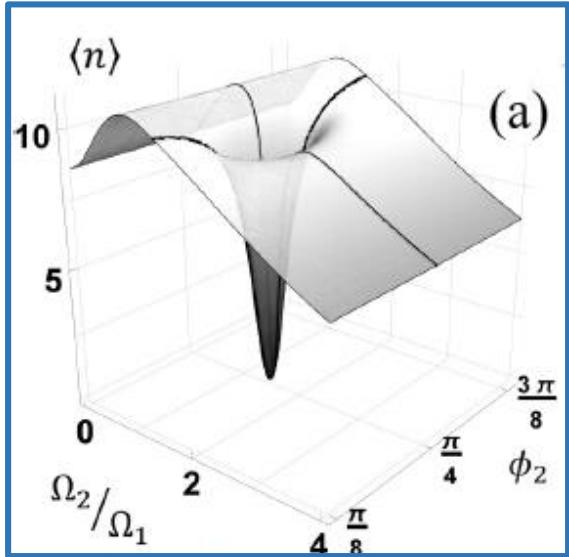
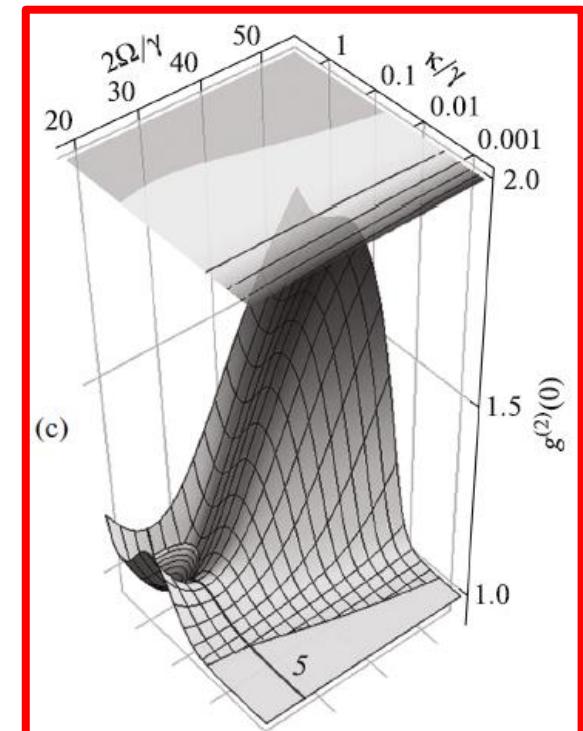
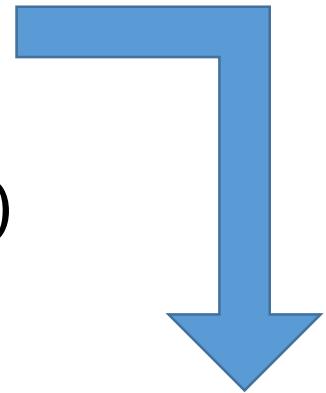


Emisie spontana + **Laser albastru defazat = GENERARE** fononi
Emisie spontana + **Laser rosu defazat = RACIRE** cuantica

• Modelul Teoretic

$$H = \hbar\omega_{qd}S_z + \hbar\omega_{ph}b^\dagger b + \hbar\Omega(S^+e^{-i\omega_L t} + S^-e^{i\omega_L t}) \\ + \hbar g S^+ S^-(b^\dagger + b),$$

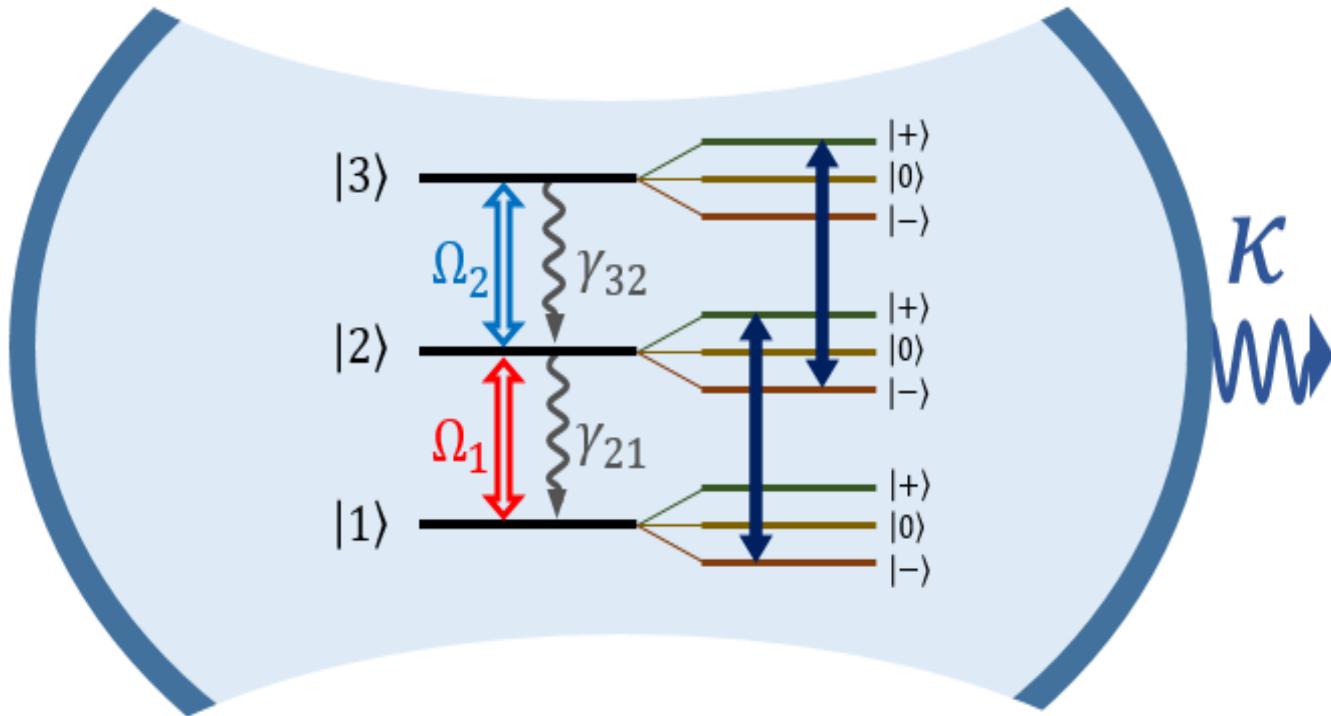
$$\dot{\rho} = -\frac{i}{\hbar}[H, \rho] + \kappa \bar{n} \mathcal{L}(b^\dagger) + \kappa(1 + \bar{n}) \mathcal{L}(b^\dagger) + \gamma \mathcal{L}(S^-) + \gamma_c \mathcal{L}(S_z)$$



$$\begin{aligned}\dot{P}_n^{(0)} &= -2|g|(P_n^{(4)} - P_n^{(3)}) + \kappa(n+1)P_{n+1}^{(0)} - \kappa n P_n^{(0)}, \\ \dot{P}_n^{(1)} &= -2|g|(P_n^{(4)} - P_n^{(3)}) + \kappa(n+1)P_{n+1}^{(1)} - (\kappa n + \alpha/2)P_n^{(1)} + \gamma_2 \cos^2 \theta P_n^{(0)}, \\ \dot{P}_n^{(2)} &= -2|g|(P_n^{(4)} + P_n^{(3)}) + \kappa(n+1)P_{n+1}^{(2)} - (\kappa n + \beta/2)P_n^{(2)}, \\ \dot{P}_n^{(3)} &= |g|n(P_{n-1}^{(1)} - P_n^{(1)} + P_{n-1}^{(2)} + P_n^{(2)})/2 - \kappa P_n^{(4)} + \kappa(n+1)P_{n+1}^{(3)} \\ &\quad - (\kappa(n-1/2) + \zeta)P_n^{(3)}, \\ \dot{P}_n^{(4)} &= |g|(n+1)(P_{n+1}^{(2)} + P_n^{(2)} - P_{n+1}^{(1)} + P_n^{(1)})/2 + \kappa(n+1)P_{n+1}^{(4)} \\ &\quad - (\kappa(n+1/2) + \zeta)P_n^{(4)},\end{aligned}$$

Punct cuantic in rezonator optic

• Punct cuantic in rezonator optic



Ce se intampla:

- Laserele pregatesc atomul intr-o superpozitie de stari
- Cavitatea optica interactioneaza cu toate tranzitiile
- Interferente cuantice se produc deoarece se interactioneaza cu mai multe tranzitii



Cu ce lucram:

- Atom cu 3 nivele
- 2 lasere
- Intuneric (vacuum electromagnetic)
- Cavitate buna
- Dipoli de tranzitie perpendiculari

Campul
electromagnetic
din cavitate
DISPARE

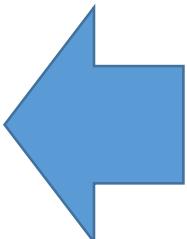
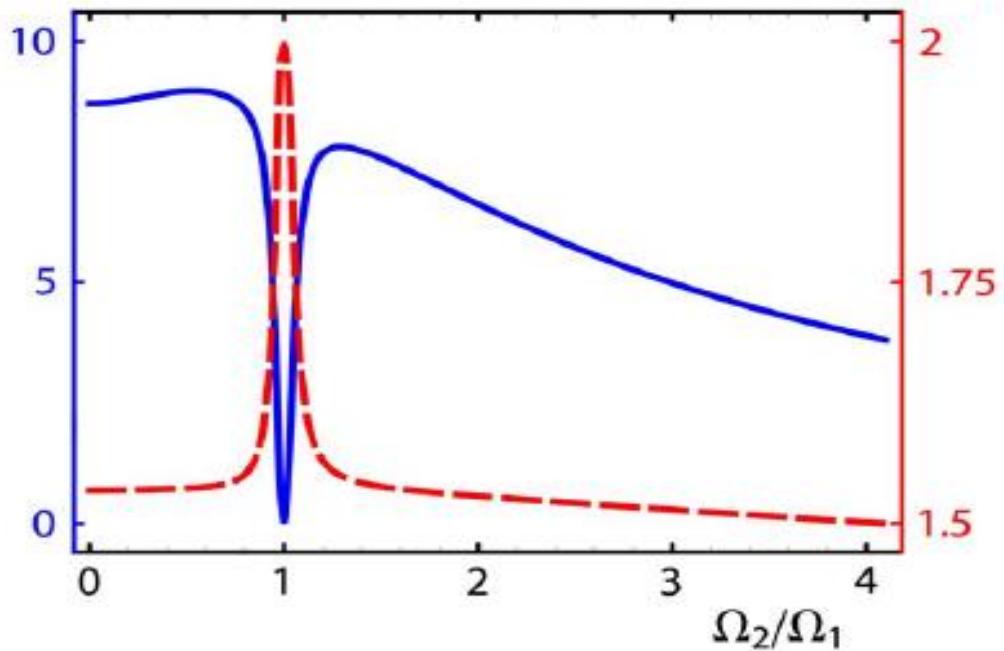
• Modelul Teoretic

$$H = \hbar\omega_c a^\dagger a + \hbar \sum_{i=1}^3 \omega_i S_{ii} + i\hbar g_1(a^\dagger S_{12} - S_{21}a) + i\hbar g_2(a^\dagger S_{23} - S_{32}a)$$

$$+ \hbar\Omega_1(S_{21}e^{-i(\omega_{L1}t+\phi_1)} + S_{12}e^{i(\omega_{L1}t+\phi_1)})$$

$$+ \hbar\Omega_2(S_{32}e^{-i(\omega_{L2}t+\phi_2)} + S_{23}e^{i(\omega_{L2}t+\phi_2)}).$$

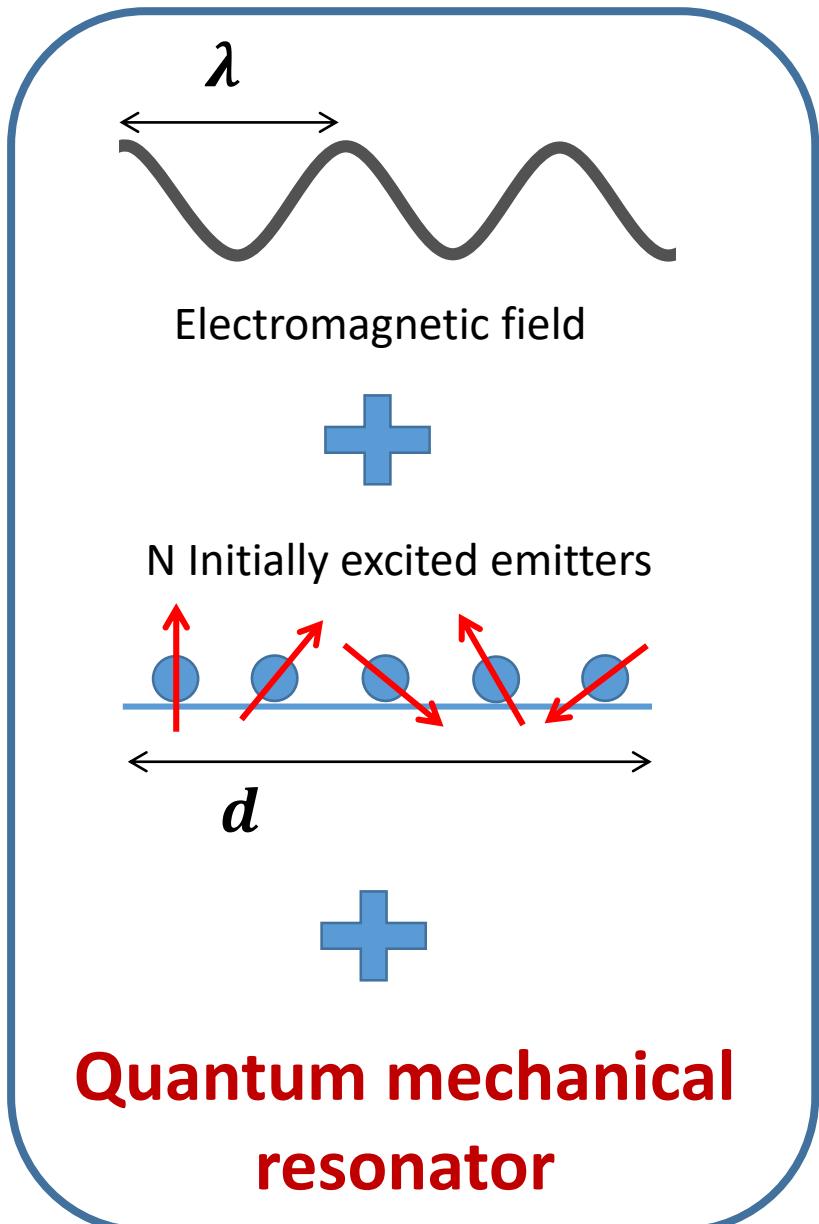
$$\frac{\partial \rho}{\partial t} = -\frac{i}{\hbar}[H, \rho] + \frac{\kappa}{2}\mathcal{L}(a) + \frac{\gamma_{32}}{2}\mathcal{L}(S_{23}) + \frac{\gamma_{21}}{2}\mathcal{L}(S_{12}).$$



$$\begin{aligned} \dot{P}_n^{(0)} &= -2g(P_n^{(4)} - P_n^{(3)}) + \kappa(n+1)P_{n+1}^{(0)} - \kappa n P_n^{(0)}, \\ \dot{P}_n^{(1)} &= -2g(P_n^{(4)} - P_n^{(3)}) + \kappa(n+1)P_{n+1}^{(1)} \\ &\quad - (\kappa n + \alpha/2)P_n^{(1)} + \gamma_{32} \cos^2 \theta P_n^{(0)}, \\ \dot{P}_n^{(2)} &= -2g(P_n^{(4)} + P_n^{(3)}) + \kappa(n+1)P_{n+1}^{(2)} - (\kappa n + \beta/2)P_n^{(2)} \\ \dot{P}_n^{(3)} &= gn(P_{n-1}^{(1)} - P_n^{(1)} + P_{n-1}^{(2)} + P_n^{(2)})/2 - \kappa P_n^{(4)} \\ &\quad + \kappa(n+1)P_{n+1}^{(3)} - (\kappa(n-1/2) + \zeta)P_n^{(3)}, \\ \dot{P}_n^{(4)} &= g(n+1)(P_{n+1}^{(2)} + P_n^{(2)} - P_{n+1}^{(1)} + P_n^{(1)})/2 \\ &\quad + \kappa(n+1)P_{n+1}^{(4)} - (\kappa(n+1/2) + \zeta)P_n^{(4)}, \end{aligned} \quad (1)$$

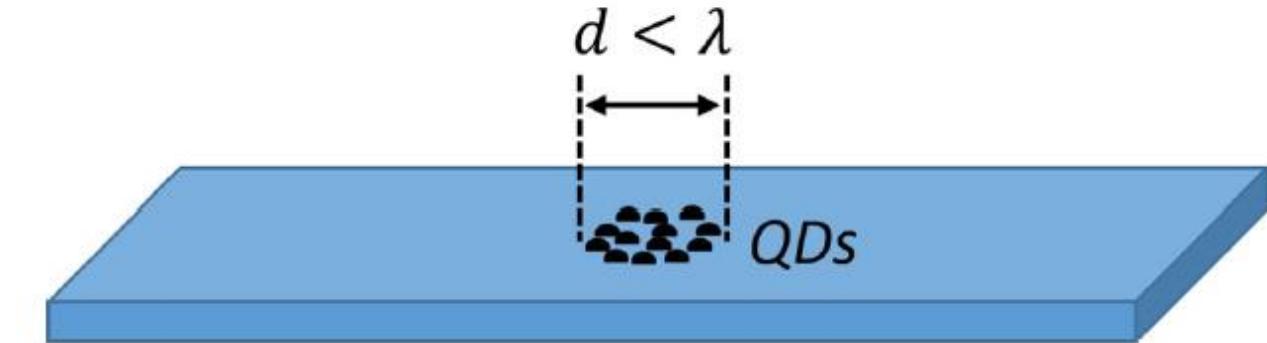
Multe puncte cuantice in rezonator mecanic

• Multe puncte cuantice in un rezonator mecanic



Cu ce lucram:

- Multe puncte cuantice aproape unul de altul
- Plasate pe un rezonator mecanic
- Excitam initial cu laser de scurta durata



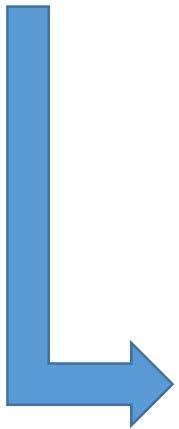
Ce se intampla:

- Deoarece sunt aproape, PC se dezexcit mai rapid si mai intens
- Interactioneaza mai rapid si mai intens cu rezonatorul mechanic
- Se produce un impuls sonor (de fononi) puternic de scurta durata

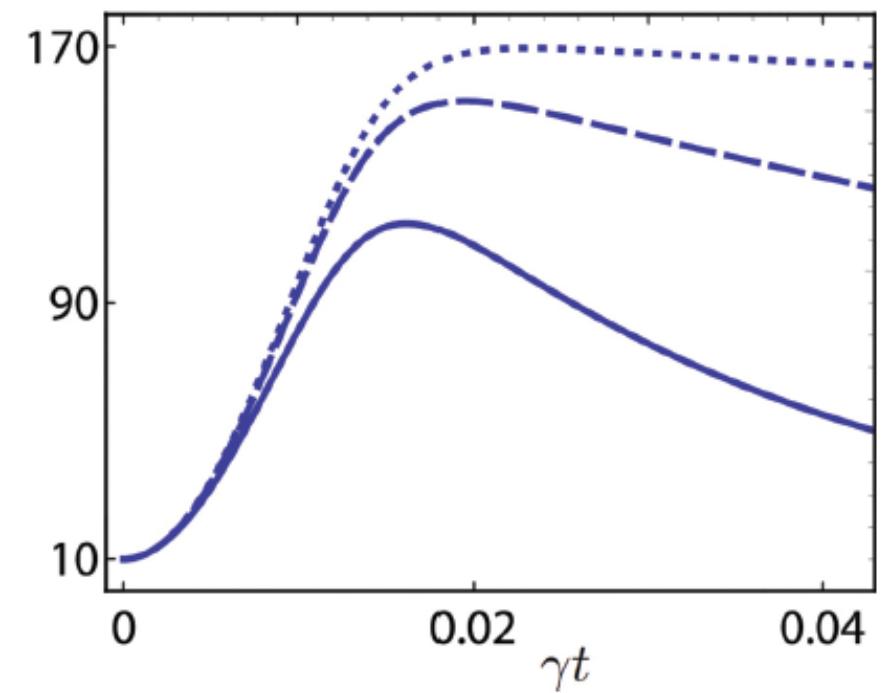
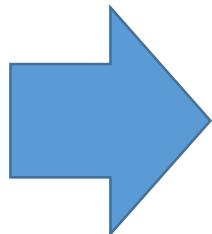
• Modelul Teoretic

$$H = \hbar\omega b^\dagger b + \hbar\omega_{\text{qd}} S_z + \hbar\eta S_{22}(b + b^\dagger).$$

$$\frac{\partial \rho}{\partial t} = -\frac{i}{\hbar}[H, \rho] + \kappa \bar{n} \mathcal{L}(b^\dagger) + \kappa(1 + \bar{n}) \mathcal{L}(b) + \gamma \mathcal{L}(S^-).$$



$$\begin{aligned}\frac{\partial \langle b^\dagger b \rangle}{\partial t} &= i\eta\{\langle S_z b \rangle - \langle S_z b^\dagger \rangle + j\langle b \rangle - j\langle b^\dagger \rangle\} \\ &\quad - 2\kappa \langle b^\dagger b \rangle + 2\kappa \bar{n}, \\ \frac{\partial \langle S_z b \rangle}{\partial t} &= -(\kappa + 2\gamma + i\omega)\langle S_z b \rangle + 2\gamma \langle S_z^2 b \rangle \\ &\quad - i\eta\{\langle S_z^2 \rangle + j\langle S_z \rangle\} - 2\gamma j(j+1)\langle b \rangle, \\ \frac{\partial \langle S_z b^\dagger \rangle}{\partial t} &= -(\kappa + 2\gamma - i\omega)\langle S_z b^\dagger \rangle + 2\gamma \langle S_z^2 b^\dagger \rangle \\ &\quad + i\eta\{\langle S_z^2 \rangle + j\langle S_z \rangle\} - 2\gamma j(j+1)\langle b^\dagger \rangle, \\ \frac{\partial \langle b \rangle}{\partial t} &= -(\kappa + i\omega)\langle b \rangle - i\eta\{\langle S_z \rangle + j\}, \\ \frac{\partial \langle b^\dagger \rangle}{\partial t} &= -(\kappa - i\omega)\langle b^\dagger \rangle + i\eta\{\langle S_z \rangle + j\}, \\ \frac{\partial \langle S_z \rangle}{\partial t} &= -2\gamma\{\langle S_z \rangle - \langle S_z^2 \rangle + j(j+1)\}.\end{aligned}$$



• Concluzii

- Studii de relevanta mondiala => diploma recunoscuta international
- Studiu teoretic nu necesita utilaj costisitor => pe larg accesibil
- Cea mai grea teorie a opticii => putin accesibila experimentatorilor

Multumesc!